

Etude d'un oscillateur linéaire amorti à un degré de liberté

1.1 Résolution avec un schéma d'EULER explicite

On le fait sous forme matricielle:

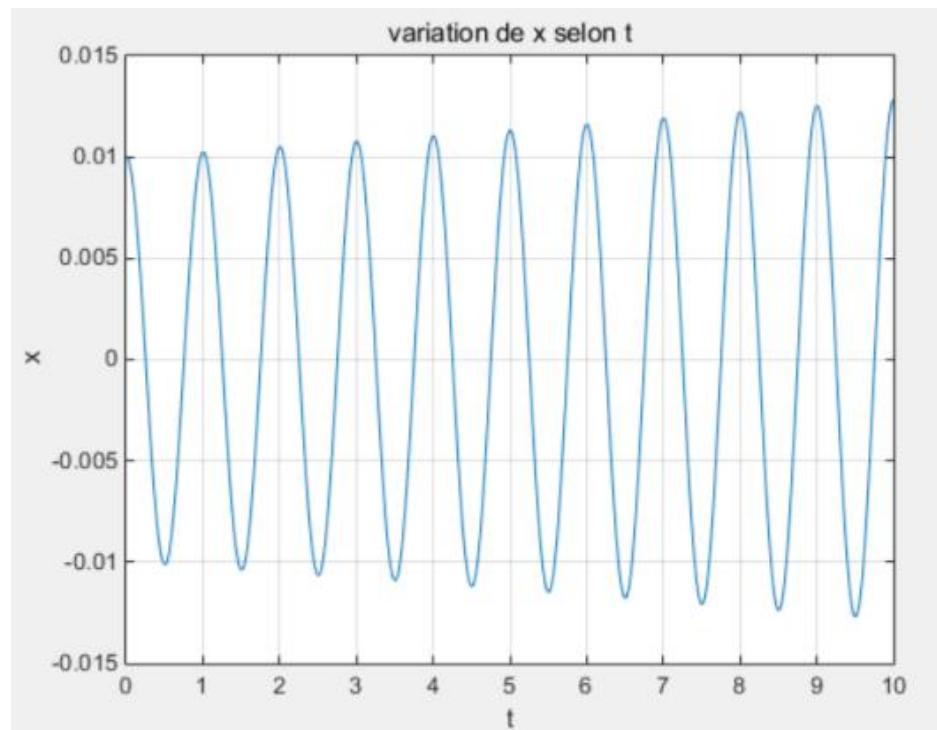
$$\begin{aligned} \begin{cases} q_{j+1} \\ \dot{q}_{j+1} \end{cases} &= \begin{cases} q_j + \Delta t \\ \dot{q}_j \end{cases} = \begin{cases} q_j + \Delta t \\ \dot{q}_j \end{cases} - \underbrace{2\zeta\omega_0\dot{q}_j - \omega_0^2 q_j}_{\text{partie amortie}} \\ &= \begin{pmatrix} 1 & \Delta t \\ -\omega_0^2 \Delta t & 1 - 2\zeta\omega_0 \Delta t \end{pmatrix} \begin{pmatrix} q_j \\ \dot{q}_j \end{pmatrix} \text{ donc } \begin{cases} q_{j+1} \\ \dot{q}_{j+1} \end{cases} = A \begin{pmatrix} q_j \\ \dot{q}_j \end{pmatrix} \end{aligned}$$

avec $A = \begin{pmatrix} 1 & \Delta t \\ -\omega_0^2 \Delta t & 1 - 2\zeta\omega_0 \Delta t \end{pmatrix}$

%% Resolution avec un schema d'Euler explicite

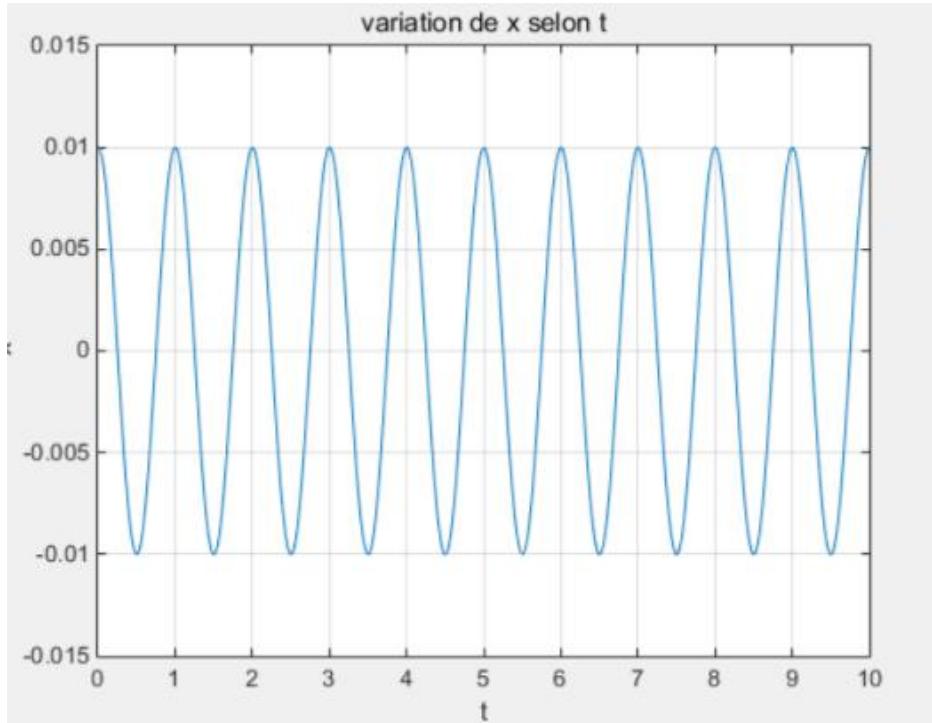
```
Aex = [1 dt; -w0^2*dt 1-2*dt*eps*w0];
```

A) Si $\Delta t > 2 * \text{eps} / \omega_0$



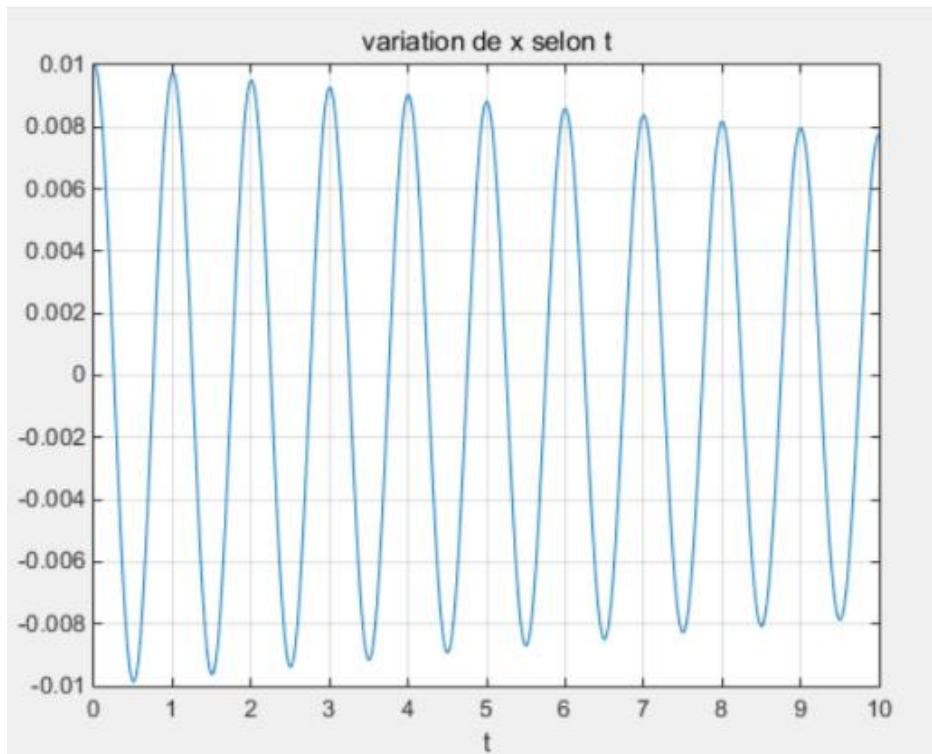
Il est instable

B) Si $dt = 2 * \text{eps}/w_0$



Il est de régime critique

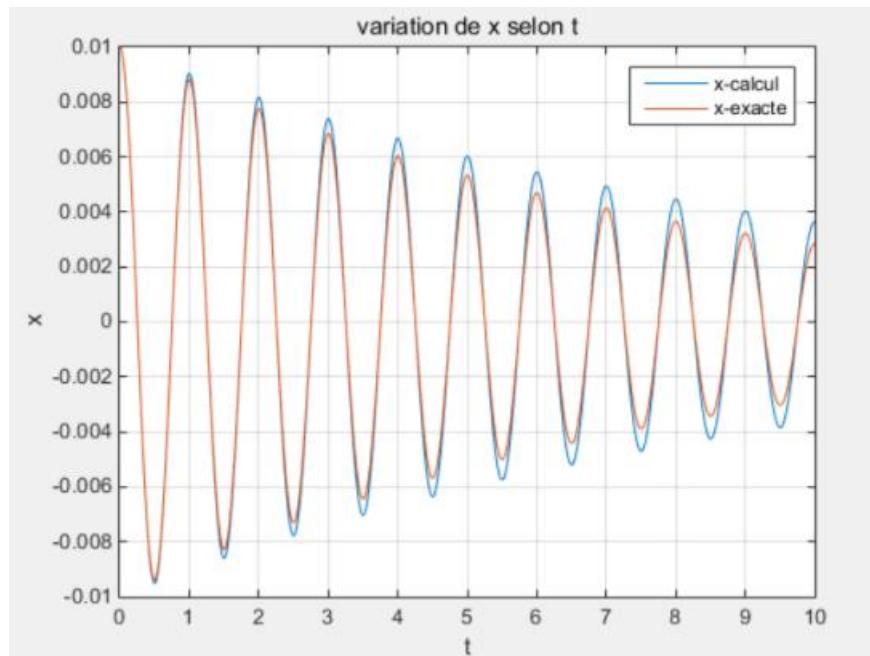
C) Si $dt = 2 * \text{eps}/w_0 * 0.8$



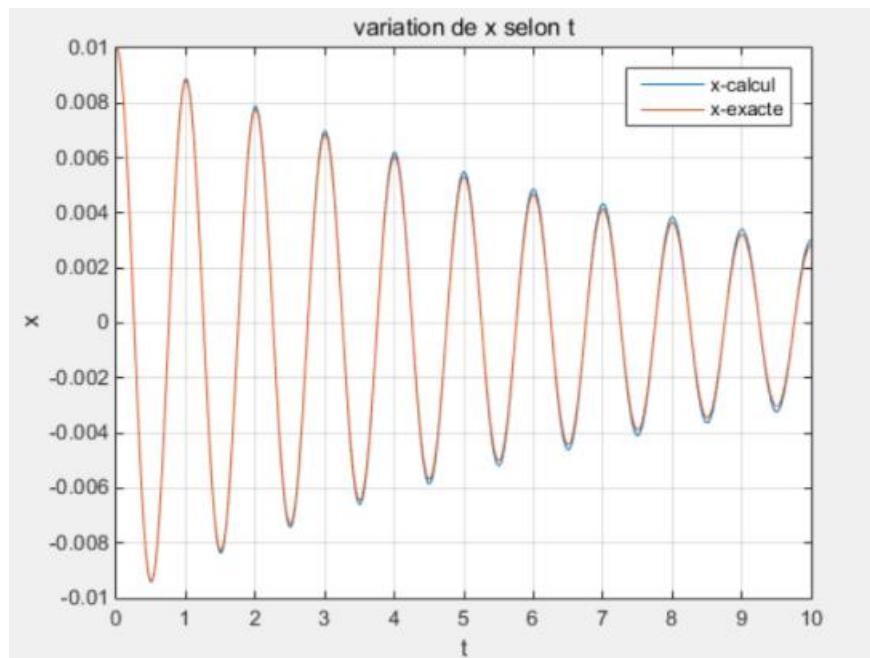
Il a tendance de converger

D)

$$dt = 0.2 * 2 * \text{eps}/w_0$$



$$dt = 0.05 * 2 * \text{eps}/w_0$$



Critère de la précision: x doit converger comme x-exact, et puis, la différence entre x-calcul et x-exacte soit être petite.

Code de 1.1

```
%% Etude de l'oscillateur lineaire amorti a un degré de liberté
clc
clear all
close all
%% Initialisation
T0 = 1;
w0 = 2*pi/T0;
eps = 0.02;
%b = 2*eps*w0*m;
x0 = 0.01;
xp0 = 0;
%dt = 2*eps/w0;
%dt = 2*eps/w0*1.2;
dt = 2*eps/w0*0.05;
U(1,1) = x0;
U(2,1) = xp0;
T = [0:dt:10*T0];
omega = w0* power(1-eps^2,1/2);
xexacte = exp(-eps*w0*T).* (x0*cos(omega*T) +
(eps*w0*x0+xp0)/omega*sin(omega*T));
%% Resolution avec un schema d'Euler explicite
Aex = [1 dt;-w0^2*dt 1-2*dt*eps*w0];
Uex = U;
%% solution explicite
for j = 0:dt:10*T0
    a = round(j/dt)-156
    Uex(:,round(j/dt)+2) = Aex * Uex(:,round(j/dt)+1);
end
%% plot
Uex(:,end) = [];
figure(1)
plot([0:dt:10*T0],Uex(1,:))
grid on
xlabel('t')
ylabel('x')
title('variation de x selon t');

figure(2)
plot([0:dt:10*T0],Uex(1,:),[0:dt:10*T0],xexacte)
grid on
```

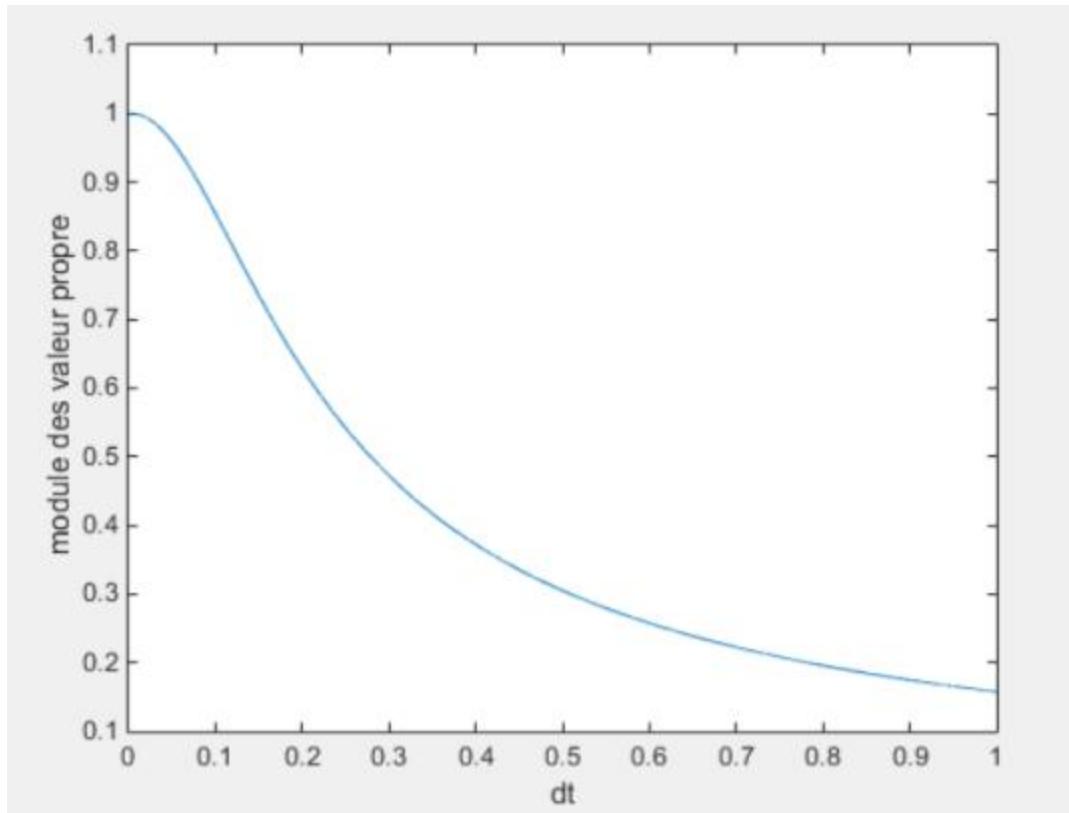
```

xlabel('t')
ylabel('x')
title('variation de x selon t');
legend('x-calcul','x-exacte')

```

1.2 Résolution avec un schéma d'EULER implicite

Si on fait avec un schéma d'Euler implicite, on calcule les module des valeur propres de la matrice d'amplification en dt différente



On a vu que la fonction de module de valeur propre de la matrice d'amplification est décroissante, donc il n'y a pas de pas de temps critique.

Code de 1.2

```

%% Implicite
Aim = inv(Aex);% determiner la matrice d'amplification Aim
dt_tt = 0:0.0001:1;
%Chercher le pas de temps critique
for i = 1:length(dt_tt)
    Aex_tt{1,i} = [1 dt_tt(i);-w0^2*dt_tt(i) 1-2*dt_tt(i)*eps*w0];

```

```

Aim_tt{1,i} = inv(Aex_tt{1,i});
vpAim_tt(:,i) = eig(Aim_tt{1,i});
module_tt(i) = sqrt(real(vpAim_tt(1,i))^2+imag(vpAim_tt(1,i))^2);
end
plot(dt_tt,module_tt);
xlabel('dt');
ylabel('module des valeur propre');

```

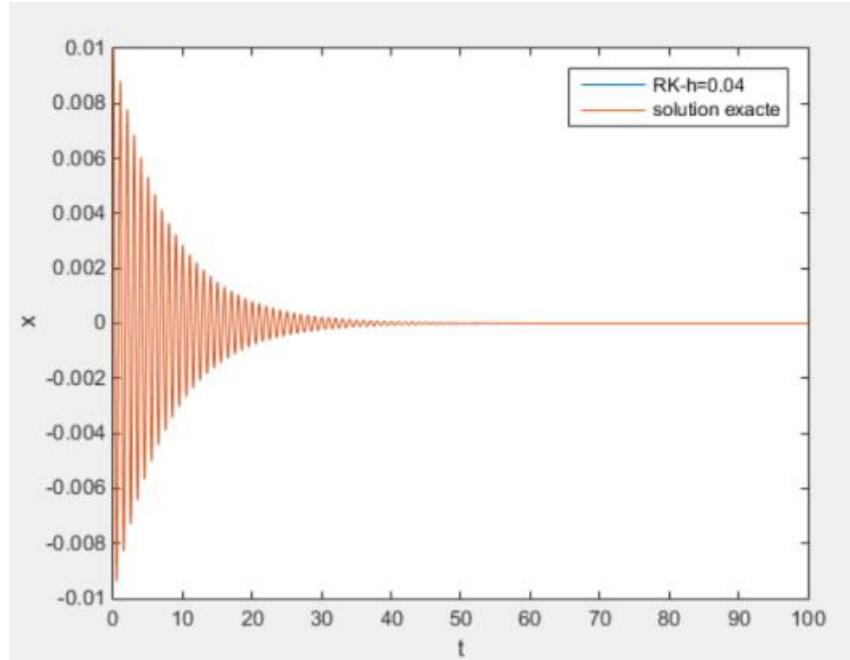
1.3a

```

h = 0.04;
%h = 0.96;
%h = 1.04;

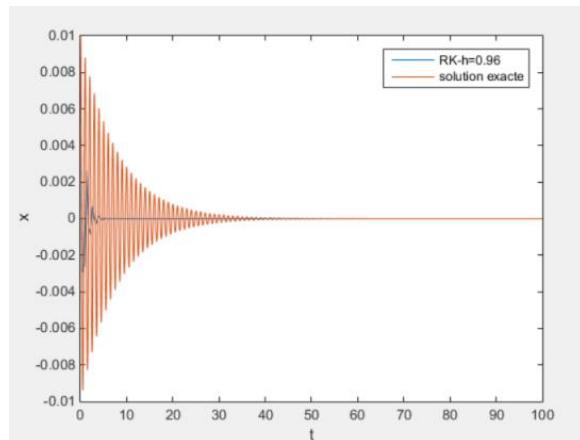
```

Quand $h = 0.04$



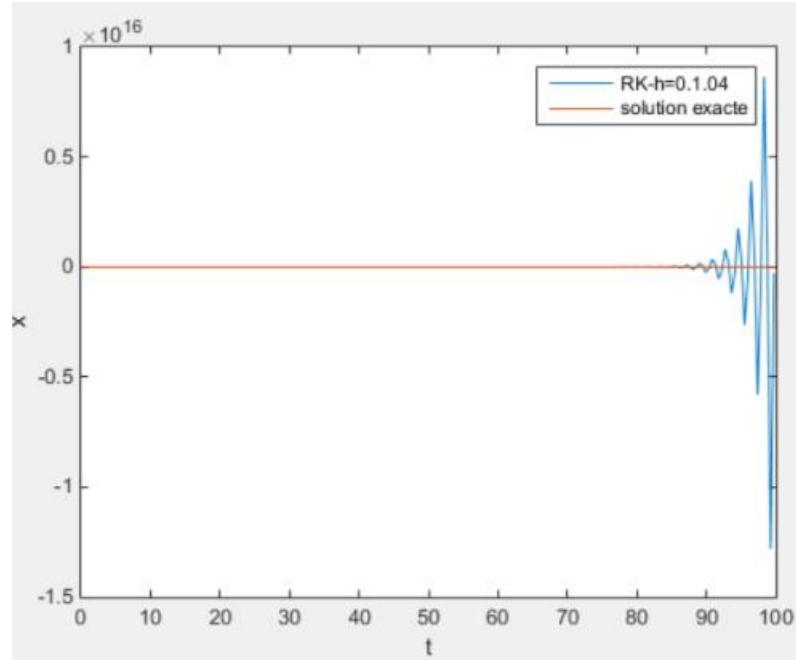
Le résultat est très proche à la solution exacte, stable et précis.

Quand $h = 0.96$



Le résultat est stable mais pas précis

Quand $h = 1.04$

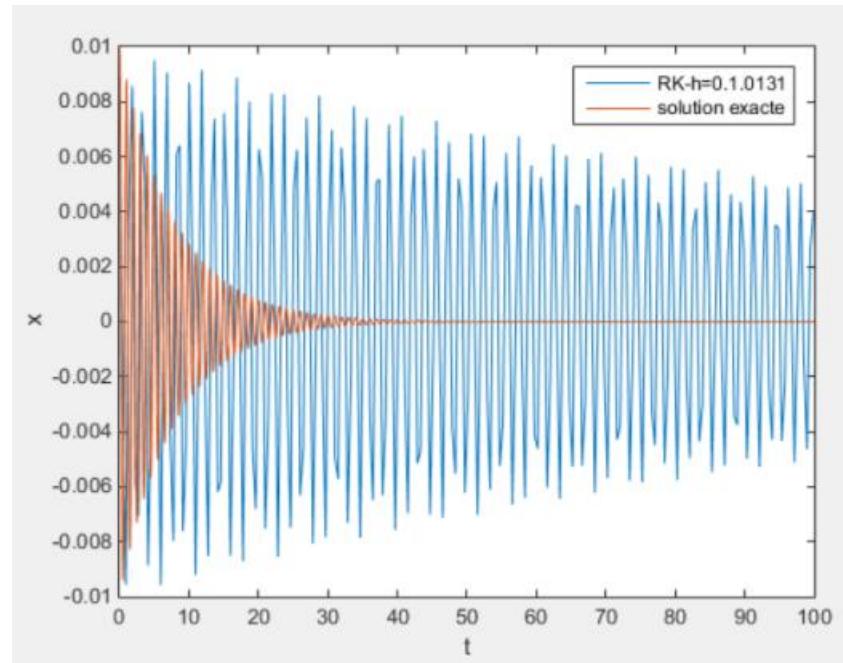


Le résultat n'est pas stable ni précis, il diverge

1.3b

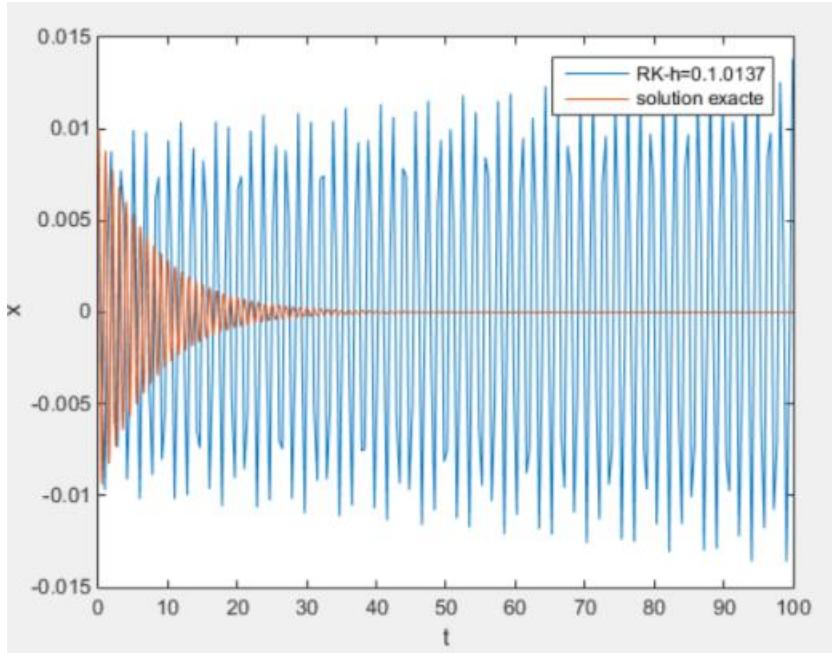
On essaye comme suivant

Quand $h = 0.10131$



Convergente

Quand $h = 1.0137$



Divergent

Donc

et on a

$$dt = h_c \times \frac{2\sqrt{2}}{\omega_0} \quad | \\ 1.0131 < hc < 1.0137$$

Code de 1.3

```
%% resolution avec un schema RUNGE KUTTA
% 1.3a
%h = 0.04;
% h = 0.96;
h = 1.0131
%h = 1.04;
dt_rk = h*2*sqrt(2)/w0;
t_rk = 0:dt_rk:100*T0;
Ark = [0,1;-w0^2,-2*eps*w0];
Urk(:,1) = [x0;xp0];
for i = 1:length(t_rk)-1
    k1 = Ark * Urk(:,i);
    k2 = Ark * (Urk(:,i) + k1*dt_rk/2);
    k3 = Ark * (Urk(:,i) + k2*dt_rk/2);
    k4 = Ark * (Urk(:,i) + k3*dt_rk);
    K = (k1+2*k2+2*k3+k4)/6;
```

```

Urk(:,i+1) = Urk(:,i) + K*dt_rk;
end
figure(2)
plot(t_rk,Urk(1,:),T,xexacte)
%legend('RK-h=0.04','solution exacte')
%legend('RK-h=0.96','solution exacte')
legend('RK-h=0.1.0131','solution exacte')
xlabel('t');
ylabel('x');

```

Etude d'un double pendule avec l'hypothèse des petits mouvements

1. Newmark explicite, gamma=0.5 beta=0

1.1 On en deduit la matrice d'amplification:

$$\begin{aligned}
\text{Preuve : } & \begin{cases} q_{n+1} = q_n + \delta t \dot{q}_n + \delta t^2 (0.5 - \beta) \ddot{q}_n + \delta t^2 \beta \ddot{q}_{n+1} \\ \dot{q}_{n+1} = \dot{q}_n + \delta t (1 - \gamma) \ddot{q}_n + \delta t \gamma \ddot{q}_{n+1} \\ \ddot{q}_{n+1} = -\omega_0^2 q_{n+1} \end{cases} \\
\Rightarrow & \underbrace{\begin{bmatrix} 1 + \beta \delta t^2 \omega_0^2 & 0 \\ \gamma \delta t \omega_0^2 & 1 \end{bmatrix}}_B \begin{pmatrix} q_{n+1} \\ \dot{q}_{n+1} \end{pmatrix} = \underbrace{\begin{bmatrix} 1 - \delta t^2 (0.5 - \beta) \omega_0^2 & \delta t \\ -(1 - \gamma) \delta t \omega_0^2 & 1 \end{bmatrix}}_C \begin{pmatrix} q_n \\ \dot{q}_n \end{pmatrix} \\
\Rightarrow & \begin{pmatrix} q_{n+1} \\ \dot{q}_{n+1} \end{pmatrix} = A \begin{pmatrix} q_n \\ \dot{q}_n \end{pmatrix} \text{ avec } A = B^{-1} \cdot C \\
\text{et } A = & \begin{bmatrix} 1 - \frac{\omega_0^2 \delta t^2}{2(1 + \beta \omega_0^2 \delta t^2)} & \frac{\delta t}{1 + \beta \omega_0^2 \delta t^2} \\ -\omega_0^2 \delta t \left[\frac{\gamma \omega_0^2 \delta t^2}{2(1 + \beta \omega_0^2 \delta t^2)} + 1 \right] & 1 - \frac{\gamma \omega_0^2 \delta t^2}{1 + \beta \omega_0^2 \delta t^2} \end{bmatrix}
\end{aligned}$$

1.2

D'apres le cour quand gamma = 0.5, il faut que :

$$\left(\gamma + \frac{1}{2}\right)^2 - 4\beta \leq \frac{4}{\omega_0^2 \Delta t^2}$$

On a maintenant gamma = 0.5 et beta = 0, donc dt < 2/w0, alors le pas de temps critique est **0.4515**

1.3

$$m\ddot{a}^2 \begin{bmatrix} 2 & 1 \\ 1 & 1 \end{bmatrix} \ddot{q}_0 + m g a \begin{bmatrix} 2 & 0 \\ 0 & 1 \end{bmatrix} q_0 = 0$$

$$\Leftrightarrow \ddot{q}_0 = -\frac{g}{a} \begin{bmatrix} 1 & -1 \\ -1 & 2 \end{bmatrix} \begin{bmatrix} 2 & 0 \\ 0 & 1 \end{bmatrix} q_0$$

$$= -\frac{g}{a} \begin{bmatrix} 2 & -1 \\ -2 & 2 \end{bmatrix} q_0$$

1.4

Si on ajoute la force qui est dépendante du temps

$$q_{nt+1} = q_n + \alpha t \dot{q}_n + \alpha t^2 (0.5 - \beta) \ddot{q}_n + \alpha t^2 \beta \ddot{q}_{nt+1}$$

et

$$\dot{q}_{nt+1} = \dot{q}_n + \alpha t (1 - \gamma) \ddot{q}_n + \alpha t \gamma \ddot{q}_{nt+1}$$

$$\Rightarrow m\ddot{a}^2 \begin{bmatrix} 2 & 1 \\ 1 & 1 \end{bmatrix} (\ddot{q}_n + m g a \begin{bmatrix} 2 & 0 \\ 0 & 1 \end{bmatrix} q_n) = F_0 \sin(\omega t) \left| \frac{a}{\sqrt{2}} \right.$$

1.5

%% premier cas Q1

```

gamma1 = 0.5;
beta1 = 0;
dt = 0.02;
T = 0:dt:T0;
M = inv([2,1;1,1])*m*g*a*[2,0;0,1]/(m*a^2);
M1 = F0*inv([2,1;1,1])*[a;a/sqrt(2)]/(m*a^2);

B = [[1 0;0 1]+beta1*dt^2*M 0*[1 0;0 1];gamma1*dt*M [1 0;0 1]];
C = [[1 0;0 1]-(0.5-beta1)*dt^2*M dt*[1 0;0 1];-(1-gamma1)*dt*M [1 0;0
1] ];
A = inv(B)*C;

vp = eig(A)

rl = [real(vp(2));real(vp(4))];

```

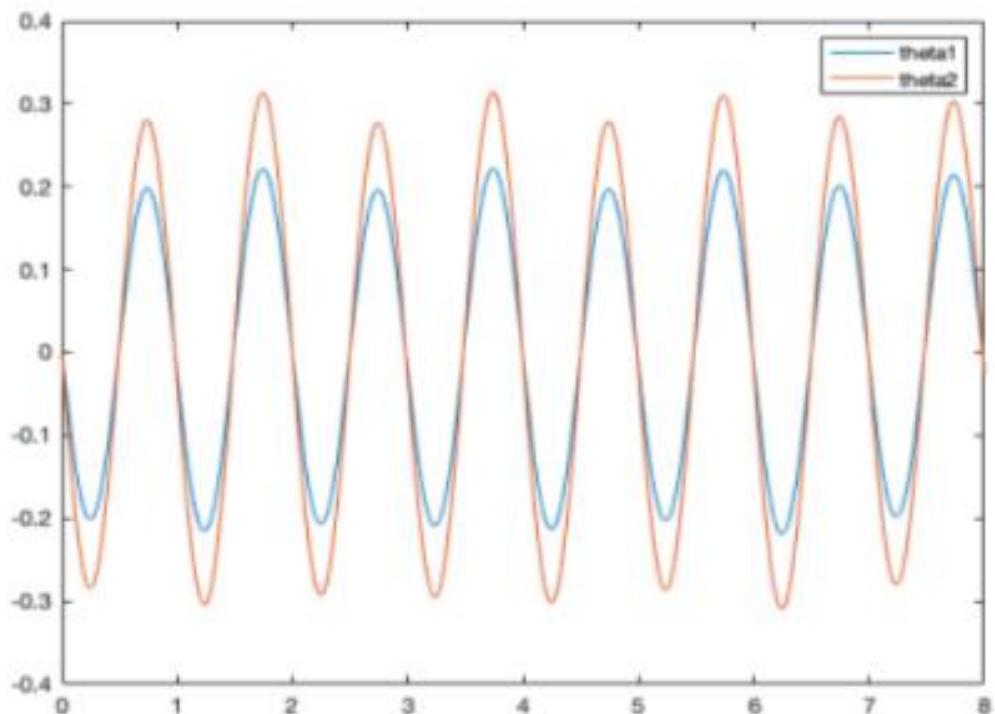
```

ig = [imag(vp(2));imag(vp(4))];
nm = [abs(vp(2));abs(vp(4))];
U(:,1) = [theta0;thetap0];
for i = 1:length(T)-1
    F =
[dt^2*(0.5-beta1)*M1*sin(w*i*dt);dt*(1-gamma1)*M1*sin(w*i*dt)];
    U(:,i+1) = A*U(:,i)+inv(B)*F;
end

plot(T,U(1,:),T,U(2,:))
grid on
legend('theta1','theta2')

```

On obtient le résultat comme suivant



1.6 S

Quand $t = 0$

<pre> theta = </pre>	<pre> thetapp = </pre>
<pre> 0 0 </pre>	<pre> -1.3152 -1.8600 </pre>
	<pre> 1.0e-14 * </pre>
	<pre> 0.0717 0.1015 </pre>

Quand $t = dt$

theta =
-0.0262
-0.0370

thetap =
-1.3048
-1.8453

thetapp =
0.3006
0.4251

Quand $t = 2dt$

theta =
-0.0519
-0.0734

thetap =
-1.2813
-1.8120

thetapp =
0.5965
0.8436

Quand $t = 0.5$

theta =
-0.1817
-0.2569

thetap =
0.6934
0.9806

thetapp =
2.0879
2.9527

2. Newmark implicite, gamma=0.5 beta=0.25

2.1 De même avec partie 1

Preuve : $\begin{cases} q_{n+1} = q_n + \delta t \dot{q}_n + \delta t^2 (0.5 - \beta) \ddot{q}_n + \delta t^2 \beta \ddot{q}_{n+1} \\ \dot{q}_{n+1} = \dot{q}_n + \delta t (1-\gamma) \ddot{q}_n + \delta t \gamma \ddot{q}_{n+1} \\ \ddot{q}_{n+1} = -\omega_0^2 q_{n+1} \end{cases}$

$$\Rightarrow \underbrace{\begin{bmatrix} 1 + \beta \delta t^2 \omega_0^2 & 0 \\ \gamma \delta t \omega_0^2 & 1 \end{bmatrix}}_B \begin{pmatrix} q_{n+1} \\ \dot{q}_{n+1} \end{pmatrix} = \underbrace{\begin{bmatrix} 1 - \delta t^2 (0.5 - \beta) \omega_0^2 & \delta t \\ - (1 - \gamma) \delta t \omega_0^2 & 1 \end{bmatrix}}_C \begin{pmatrix} q_n \\ \dot{q}_n \end{pmatrix}$$

$$\Rightarrow \begin{pmatrix} q_{n+1} \\ \dot{q}_{n+1} \end{pmatrix} = A \begin{pmatrix} q_n \\ \dot{q}_n \end{pmatrix} \text{ avec } A = B^{-1} \cdot C$$

et $A = \left[\begin{array}{cc} \frac{\omega_0^2 \delta t^2}{1 - \frac{\delta t^2}{2(1 + \beta \omega_0^2 \delta t^2)}} & \frac{\delta t}{1 + \beta \omega_0^2 \delta t^2} \\ -\omega_0^2 \delta t \left[\frac{\gamma \omega_0^2 \delta t^2}{2(1 + \beta \omega_0^2 \delta t^2)} + 1 \right] & 1 - \frac{\gamma \omega_0^2 \delta t^2}{1 + \beta \omega_0^2 \delta t^2} \end{array} \right]$

2.2

A partir de partie 1, il faut juste changer la valeur de beta = 0.25

2.3

$$ma^2 \begin{bmatrix} 2 & 1 \\ 1 & 1 \end{bmatrix} \ddot{q}_0 + mga \begin{bmatrix} 2 & 0 \\ 0 & 1 \end{bmatrix} q_0 = 0$$

$$\Leftrightarrow \ddot{q}_0 = -\frac{g}{a} \begin{bmatrix} 1 & -1 \\ -1 & 2 \end{bmatrix} \begin{bmatrix} 2 & 0 \\ 0 & 1 \end{bmatrix} q_0$$

$$= -\frac{g}{a} \begin{bmatrix} 2 & -1 \\ -2 & 2 \end{bmatrix} q_0$$

2.4

Si on ajoute la force qui est dépendante du temps

$$q_{n+1} = q_n + \alpha t \dot{q}_n + \alpha t^2 (0.5 - \beta) \ddot{q}_n + \alpha t^2 \beta \ddot{q}_{n+1}$$

et

$$\dot{q}_{n+1} = \dot{q}_n + \alpha t (1 - \gamma) \ddot{q}_n + \alpha t \gamma \ddot{q}_{n+1}$$

$$\Rightarrow ma^2 \begin{bmatrix} 2 & 1 \\ 1 & 1 \end{bmatrix} (\ddot{q}_n + mga \begin{bmatrix} 2 & 0 \\ 0 & 1 \end{bmatrix} q_n) = F_0 \sin(\omega t) \left| \frac{a}{\sqrt{2}} \right.$$

2.5

%% premier cas Q1

```

gamma1 = 0.5;
beta1 = 0.25;
dt = 0.02;
T = 0:dt:T0;
M = inv([2,1;1,1])*m*g*a*[2,0;0,1]/(m*a^2);
M1 = F0*inv([2,1;1,1])*[a;a/sqrt(2)]/(m*a^2);

B = [[1 0;0 1]+beta1*dt^2*M 0*[1 0;0 1];gamma1*dt*M [1 0;0 1]];
C = [[1 0;0 1]-(0.5-beta1)*dt^2*M dt*[1 0;0 1];-(1-gamma1)*dt*M [1 0;0
1]];
A = inv(B)*C;

```

```

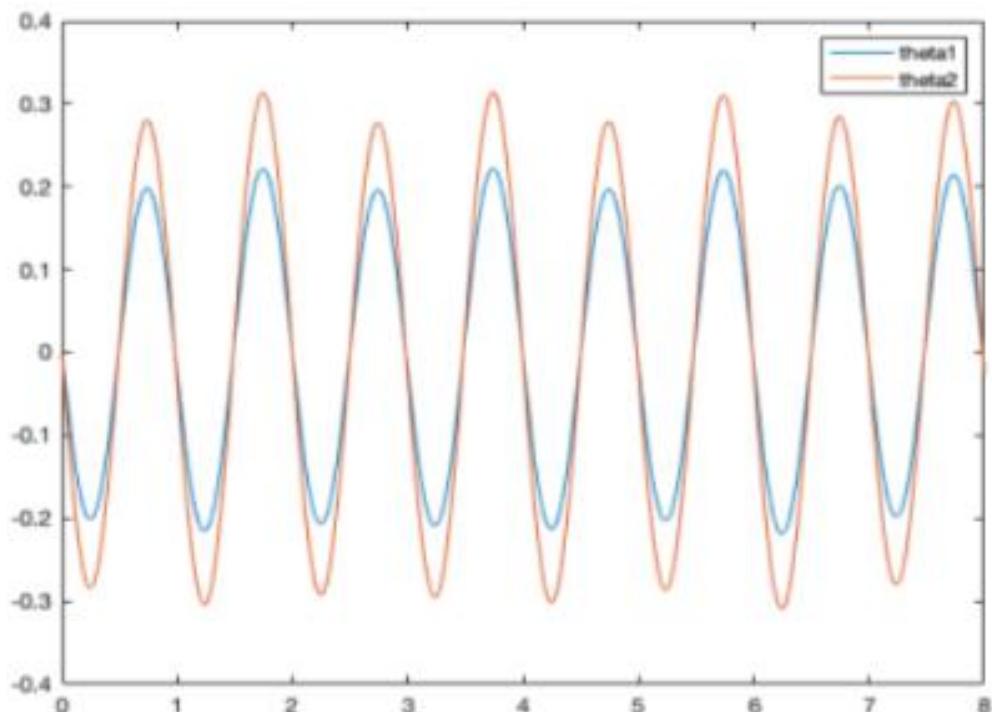
vp = eig(A)

rl = [real(vp(2));real(vp(4))];
ig = [imag(vp(2));imag(vp(4))];
nm = [abs(vp(2));abs(vp(4))];
U(:,1) = [theta0;thetap0];
for i = 1:length(T)-1
    F =
[dt^2*(0.5-beta1)*M1*sin(w*i*dt);dt*(1-gamma1)*M1*sin(w*i*dt)];
    U(:,i+1) = A*U(:,i)+inv(B)*F;
end

plot(T,U(1,:),T,U(2,:))
grid on
legend('thetal1','thetal2')

```

On obtient le resultat comme suivant



Quand $t = 0$

<pre> theta = 0 0 </pre>	<pre> thetap = -1.3152 -1.8600 0.0717 0.1015 </pre>	<pre> thetapp = 1.0e-14 * </pre>
--------------------------	---	----------------------------------

Quand t = dt

```
theta =  
-0.0262  
-0.0370
```

```
thetap =  
-1.3048  
-1.8453
```

```
thetapp =  
0.3006  
0.4251
```

Quand t = 2dt

```
theta =  
-0.0519  
-0.0734
```

```
thetap =  
-1.2813  
-1.8120
```

```
thetapp =  
0.5965  
0.8436
```

Quand t = 0.5

```
theta =  
-0.1817  
-0.2569
```

```
thetap =  
0.6934  
0.9806
```

```
thetapp =  
2.0879  
2.9527
```

Etude d'un oscillateur non linéaire à un degré de Liberté

1. NEWMARK explicite

1.1

$$\begin{cases} q_{j+1} = q_j + \alpha t \dot{q}_j + \alpha t^2 (0.5 - \beta) \ddot{q}_j + \beta \alpha t^2 \ddot{q}_{j+1} \\ \dot{q}_{j+1} = \dot{q}_j + \alpha t (1 - \gamma) \ddot{q}_j + \gamma \alpha t \ddot{q}_{j+1} \\ \ddot{q}_j + w_0^2 q_j (1 + a q_j^2) = 0 \end{cases}$$

1.2

Code de 1.2

```
%% Oscillateur non linéaire à 1 DDL
clc
clear all
close all

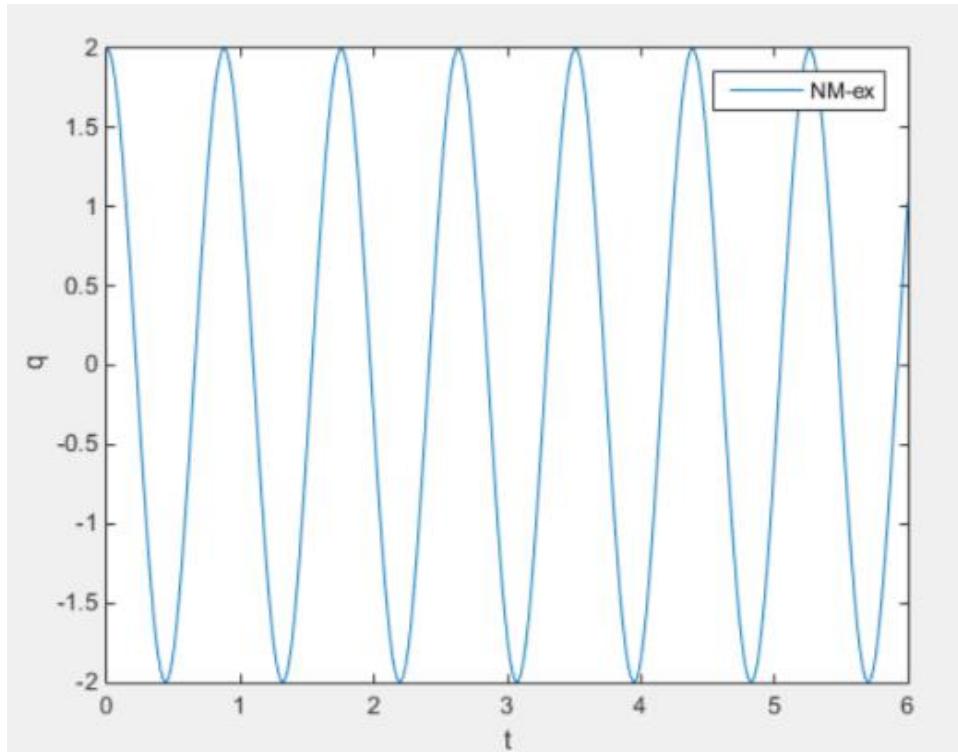
%% Initialisation global
q0 = 2;
qp0 = 0;
q(1) = q0;
qp(1) = qp0;
w0 = 2*pi;
a = 0.1;
T0 = 6;

%% NEWMARK explicite
gamma1 = 0.5;
beta1 = 0;
dt = 0.02
T = 0:dt:T0;
for j = 1:length(T) - 1
    qpp(j) = -w0^2*q(j)*(1+a*q(j)*q(j));
    q(j+1) = q(j)+dt*qp(j)+dt^2*(0.5-beta1)*qpp(j);
    qp(j+1) =
        qp(j)+dt*(1-gamma1)*qpp(j)+gamma1*dt*(-w0^2*q(j+1)*(1+a*q(j+1)*q(j+1)));
end
```

```

plot(T,q);
legend('NM-ex')
xlabel('t');
ylabel('q');

```



1.3

Code de 1.3

```
Ans1 = [q(1) q(2) q(3) q(length(T))]
```

```

Ans1 =
2.0000    1.9779    1.9123    1.0329

```

2. NEWMARK implicite

2.1

```

%% NEWMARK implicite
gamma2 = 0.5;
beta2 = 0.25;

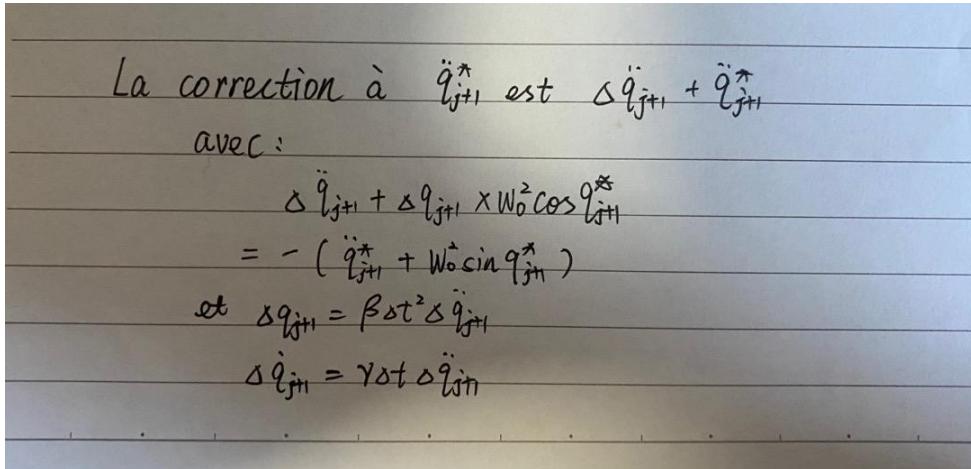
```

On cherche à minimiser la valeur absolue de $q_{pp} + w_0^2 q^2 (1 + a q^2)$

$$\ddot{q} + \omega_0^2 q (1 + a q^2)$$

Pour que les résultat soient plus proche de solution réel

2.2



2.3

Code de 2.3

```
%% NEWMARK implicite
gamma2 = 0.5;
beta2 = 0.25;

%si on choisit une precision de 0.01
d = 0.01;
qe(1) = q(1); %estimation de q
qpe(1) = qp(1); %estimation de qp
qppe(1) = qpp(1); %estimation de qpp

for i = 1:length(T)-1
    qe(i+1) = q(i)+dt*qp(i)+dt^2*(0.5-beta2)*qpp(i);
    qpe(i+1) = qp(i)+dt*(1-gamma2)*qpp(i);
    qppe(i+1) = 0;

    h = qppe(i+1) + w0^2*qe(i+1)*(1+a*qe(i+1)*qe(i+1)); %fonction de estimation
    while abs(h)>=d
        delta_qppe = -(qppe(i+1)
        + w0^2*qe(i+1)*(1+a*qe(i+1)*qe(i+1)))/(1+w0^2*beta2*dt^2*(1+3*a*qe(i+1)
        *qe(i+1)));
    end
end
```

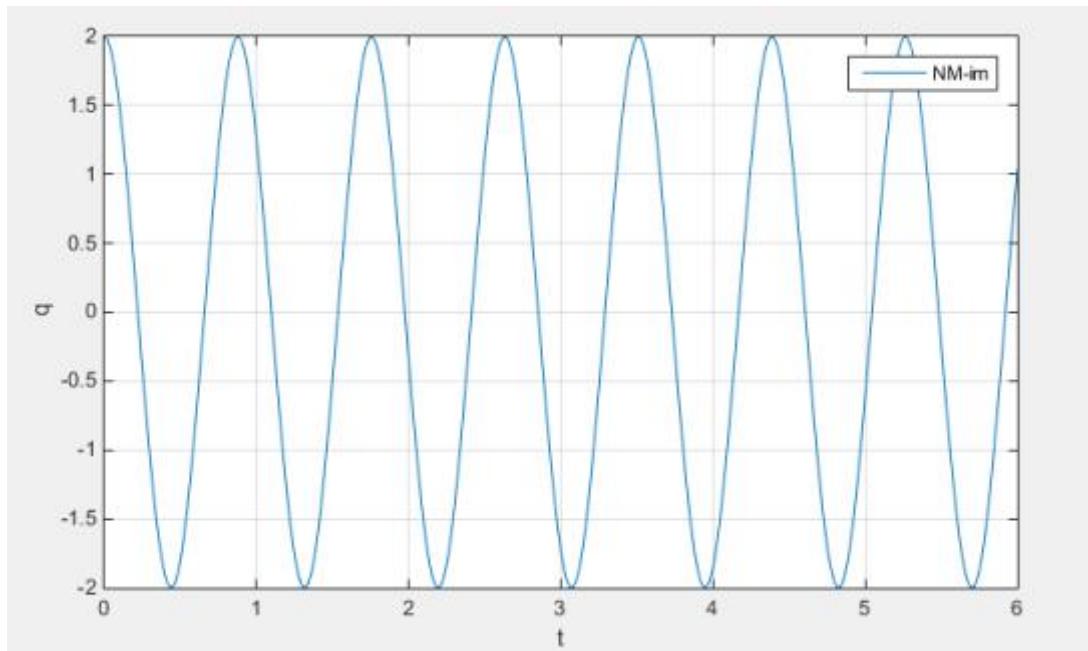
```

delta_qe = beta2*dt^2*delta_qppe;
delta_qpe = gamma2*dt*delta_qppe;

%correction
qe(i+1) = qe(i+1)+delta_qe;
qpe(i+1) = qpe(i+1)+delta_qpe;
qppe(i+1) = qppe(i+1)+delta_qppe;
h = qppe(i+1) +w0^2*qe(i+1)*(1+a*qe(i+1)*qe(i+1));
end
q(i+1) = qe(i+1);
qp(i+1) = qpe(i+1);
qpp(i+1) = qppe(i+1);
end
figure(2)
plot(T,q) ;
legend('NM-im')
xlabel('t');
ylabel('q');
grid on

```

On obtient le résultat;



2.4 on print le résultat

```
Ans2 = [q(1) q(2) q(3) q(length(t)) ]
```

|

```
Ans2 =
```

```
2.0000    1.9781    1.9131    0.8478|
```

3. énergie Mécanique

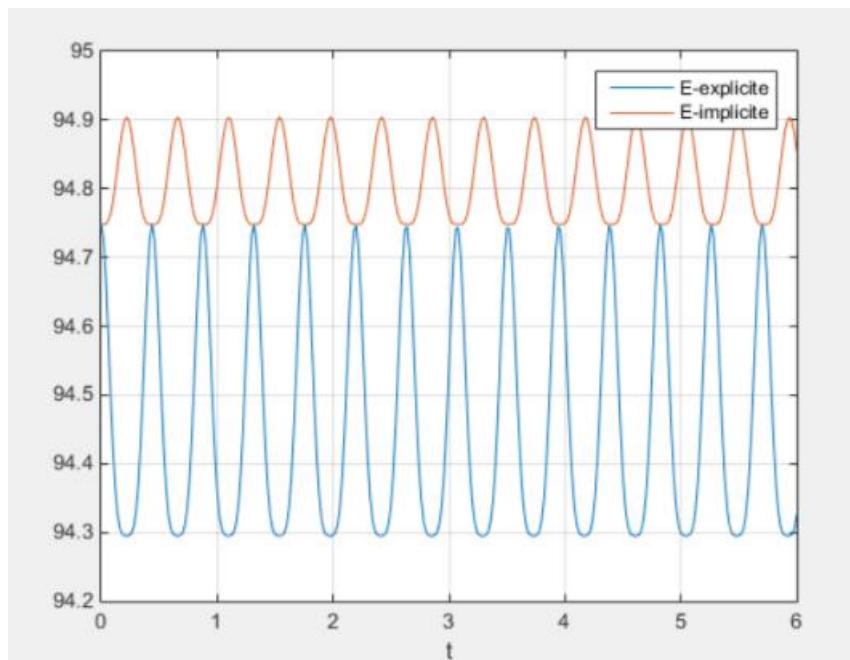
3.1

L'énergie mécanique pour cet oscillateur non linéaire est définie comme la somme de l'énergie cinétique et l'énergie potentiel.

```
%pour l'energie cinetique, c'est 0.5*dq^2  
%pour l'energie potentiel, on fait un integrale,  
%c'est 0.5*w0*w0*q*q+0.25*alpha*w0*w0*q^4
```

3.2

On compare les deux énergies obtenues par différentes intégrations:



Code de 3.2

```
%% Energie mécanique  
for n = 1:length(T)
```

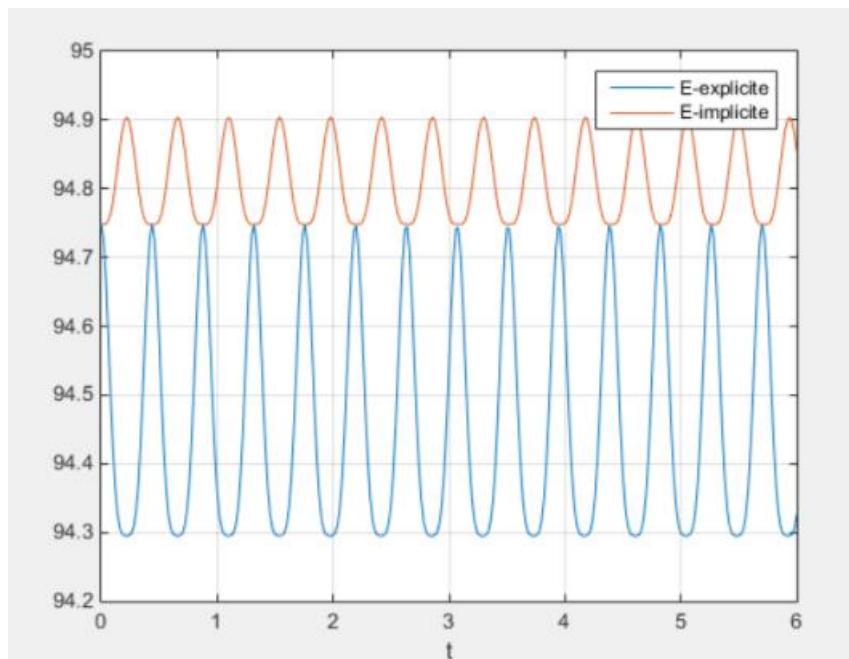
```

%explicite
Ec_ex(n) = 0.5*qp_ex(n)^2;
Ep_ex(n) = 0.5*w0^2*q_ex(n)^2+0.25*w0^2*a*q_ex(n)^4;
Etot_ex(n) = Ec_ex(n)+Ep_ex(n);
%implicite
Ec_im(n) = 0.5*qp(n)^2;
Ep_im(n) = 0.5*w0^2*q(n)^2+0.25*w0^2*a*q(n)^4;
Etot_im(n) = Ec_im(n)+Ep_im(n);
end
figure(3)
plot(T,Etot_ex,T,Etot_im)
grid on
xlabel('t')
legend('E-explicite','E-implicite')

```

3.3

Comme on a pris $dt = 0.02$, on peut conclure directement selon le résultat de question 3.2



l'énergie implicite est toujours plus grande de l'énergie explicite mais, quelque fois, ils ont la même l'énergie, et l'énergie calcule par méthode implicite est moins oscillée

Tous les code

```
%% Oscillateur non lineaire a 1 DDL
clc
clear all
close all
%% Initialisation global
q0 = 2;
qp0 = 0;
q(1) = q0;
qp(1) = qp0;
w0 = 2*pi;
a = 0.1;
T0 = 6;
qpp(1) = -w0^2*q(1)*(1+a*q(1)*q(1));

%% NEWMARK explicite
gamma1 = 0.5;
beta1 = 0;
dt = 0.02
T = 0:dt:T0;
for j = 1:length(T)
    qpp(j) = -w0^2*q(j)*(1+a*q(j)*q(j));
    q(j+1) = q(j)+dt*qp(j)+dt^2*(0.5-beta1)*qpp(j);
    qp(j+1) =
        qp(j)+dt*(1-gamma1)*qpp(j)+gamma1*dt*(-w0^2*q(j+1)*(1+a*q(j+1)*q(j+1)))
    ;
end
figure(1)
q(302)=[];
plot(T,q);
hold on
legend('NM-ex')
xlabel('t');
ylabel('q');
q_ex = q;
qp_ex = qp;
qpp1_ex = qpp;

Ans1 = [q(1) q(2) q(3) q(length(T))]

%% NEWMARK implicite
gamma2 = 0.5;
beta2 = 0.25;
```

```

%si on choisit une precision de 0.0001
d = 0.01;
qe(1) = q(1);%estimation de q
qpe(1) = qp(1); %estimation de qp
qppe(1) = qpp(1);%estimation de qpp

for i = 1:length(T)-1
    qe(i+1) = q(i)+dt*qp(i)+dt^2*(0.5-beta2)*qpp(i);
    qpe(i+1) = qp(i)+dt*(1-gamma2)*qpp(i);
    qppe(i+1) = 0;

    h = qppe(i+1) +w0^2*qe(i+1)*(1+a*qe(i+1)*qe(i+1)); %fonction de
estimation
    while abs(h)>=d
        delta_qppe = -(qppe(i+1)
+w0^2*qe(i+1)*(1+a*qe(i+1)*qe(i+1)))/(1+w0^2*beta2*dt^2*(1+3*a*qe(i+1)
*qe(i+1)));
        delta_qe = beta2*dt^2*delta_qppe;
        delta_qpe = gamma2*dt*delta_qppe;

        %correction
        qe(i+1) = qe(i+1)+delta_qe;
        qpe(i+1) = qpe(i+1)+delta_qpe;
        qppe(i+1) = qppe(i+1)+delta_qppe;
        h = qppe(i+1) +w0^2*qe(i+1)*(1+a*qe(i+1)*qe(i+1));
    end
    q(i+1) = qe(i+1);
    qp(i+1) = qpe(i+1);
    qpp(i+1) = qppe(i+1);
end

figure(2)
plot(T,q) ;
legend('NM-im')
xlabel('t');
ylabel('q');
grid on

Ans2 = [q(1) q(2) q(3) q(length(T)) ]
%% Energie mecanique
for n = 1:length(T)
    %explicite
    Ec_ex(n) = 0.5*qp_ex(n)^2;
    Ep_ex(n) = 0.5*w0^2*q_ex(n)^2+0.25*w0^2*a*q_ex(n)^4;
    Etot_ex(n) = Ec_ex(n)+Ep_ex(n);

```

```

%implicite
Ec_im(n) = 0.5*qp(n)^2;
Ep_im(n) = 0.5*w0^2*q(n)^2+0.25*w0^2*a*q(n)^4;
Etot_im(n) = Ec_im(n)+Ep_im(n);

end
figure(3)
plot(T,Etot_ex,T,Etot_im)
grid on
xlabel('t')
legend('E-explicite','E-implicite')

```